
Jet and X-ray corona of Cygnus X-1

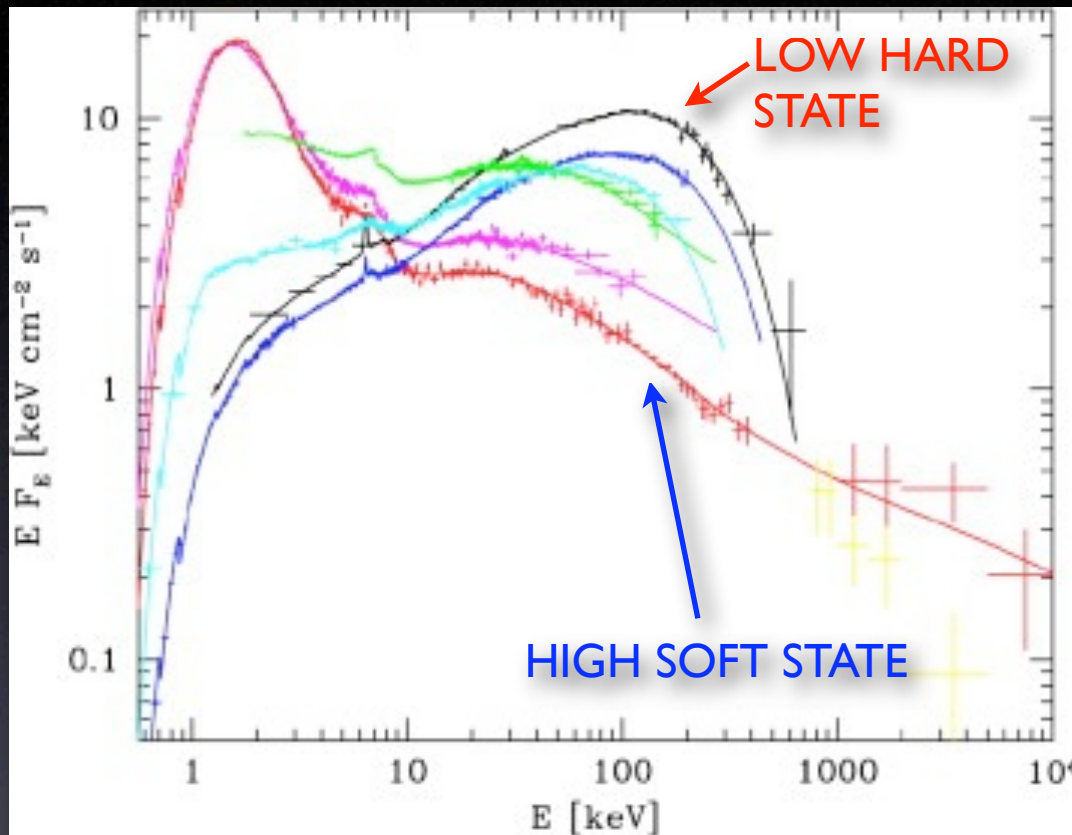
Julien Malzac (CESR/CNRS, Toulouse, France)

Outline

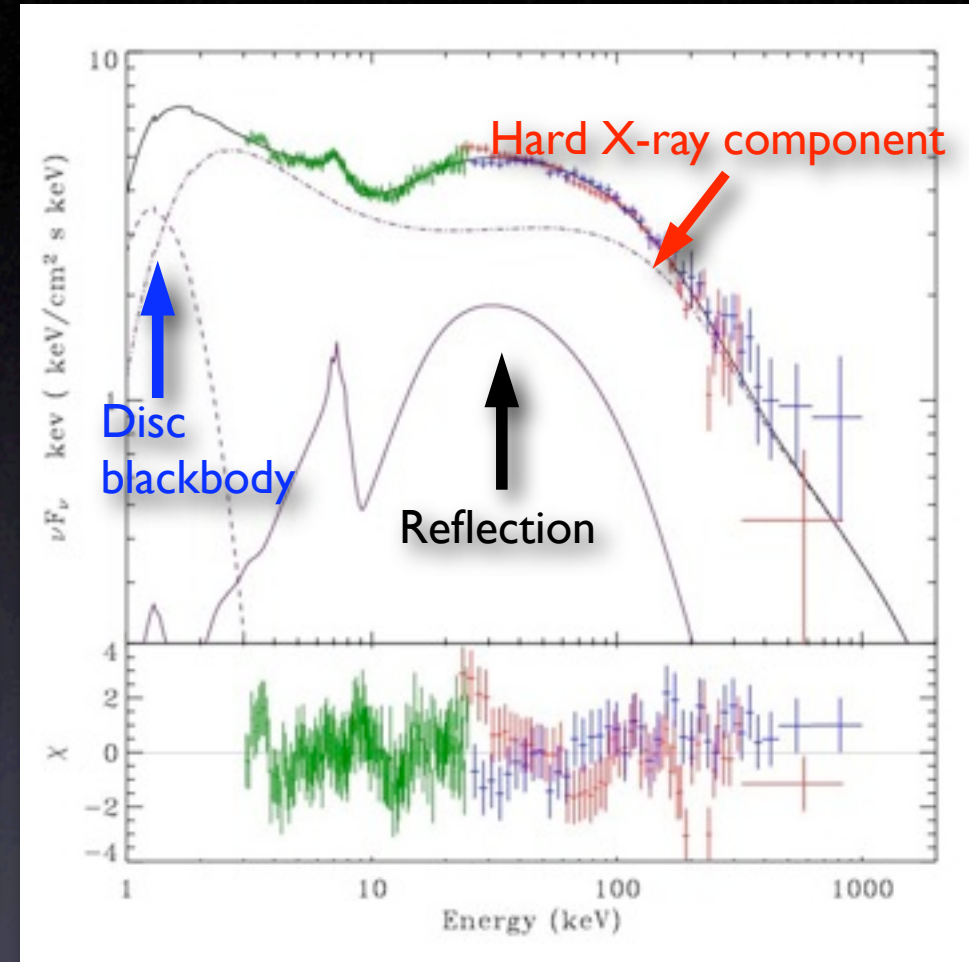
1. Black hole X-ray spectral states: observations and models
2. A numerical kinetic/radiation model for state transitions
3. Comparisons to spectra of Cygnus X-1 and GX 339-4
4. Constraints from jet power and energetics
5. TeV detection of Cygnus X-1

High energy emission of accreting black holes

Cygnus X-1



Zdziarski et al 2003



Malzac et al. 2006

LOW HARD STATE:

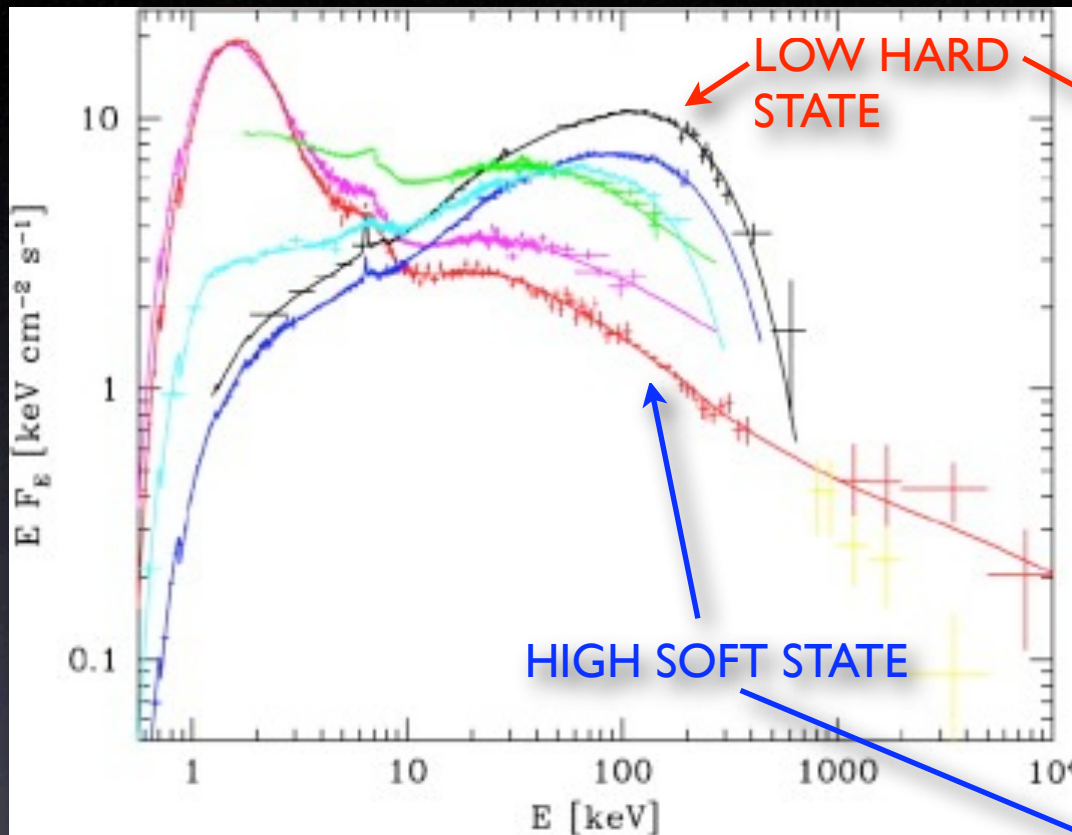
disc blackbody and reflection: weak / Corona: THERMAL Comptonisation

HIGH SOFT STATE:

disc blackbody and reflection: strong / Corona: NON-THERMAL (or hybrid) Comptonisation

High energy emission of accreting black holes

Cygnus X-1



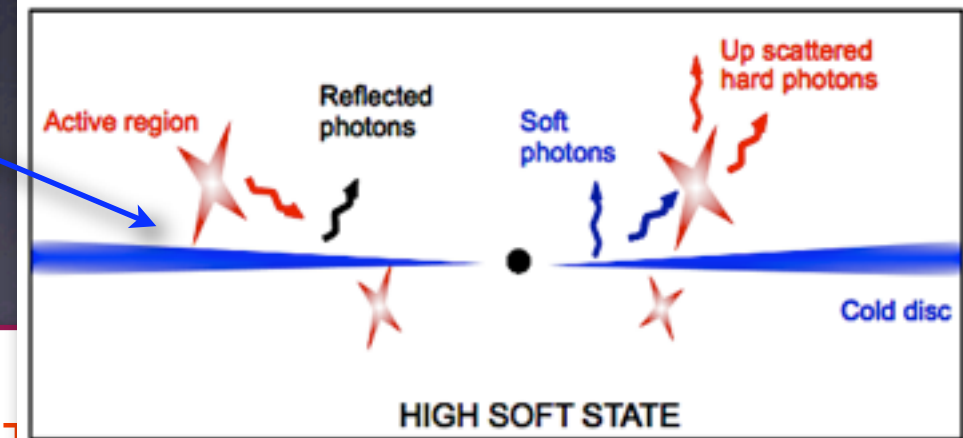
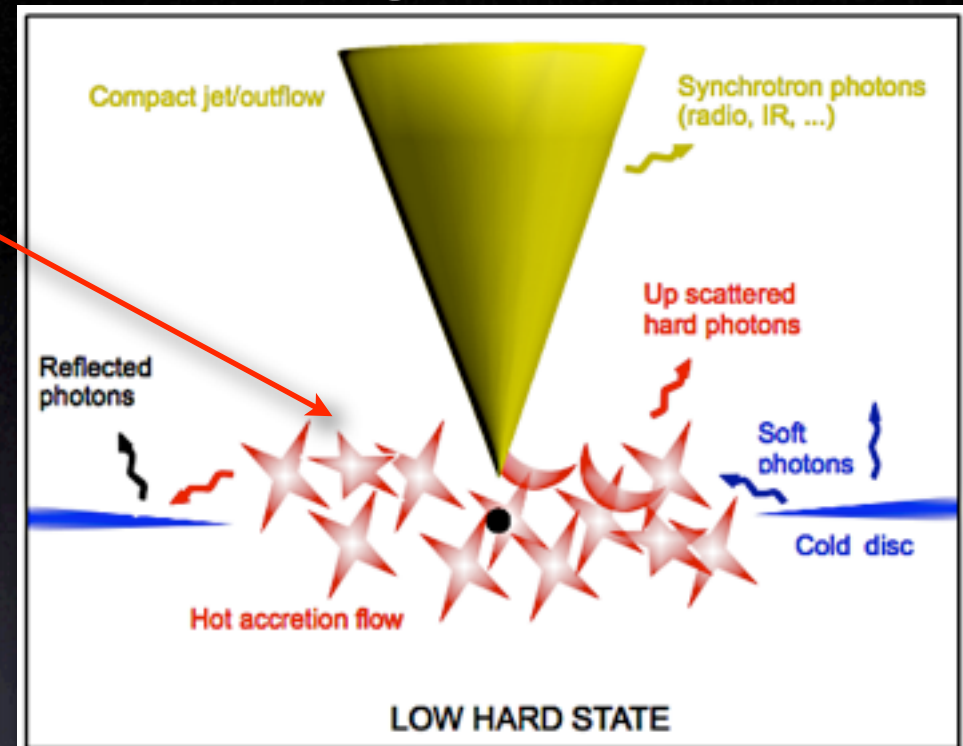
Zdziarski et al 2003

LOW HARD STATE:

disc blackbody and reflection: weak / Corona: THERMAL Comptonisation

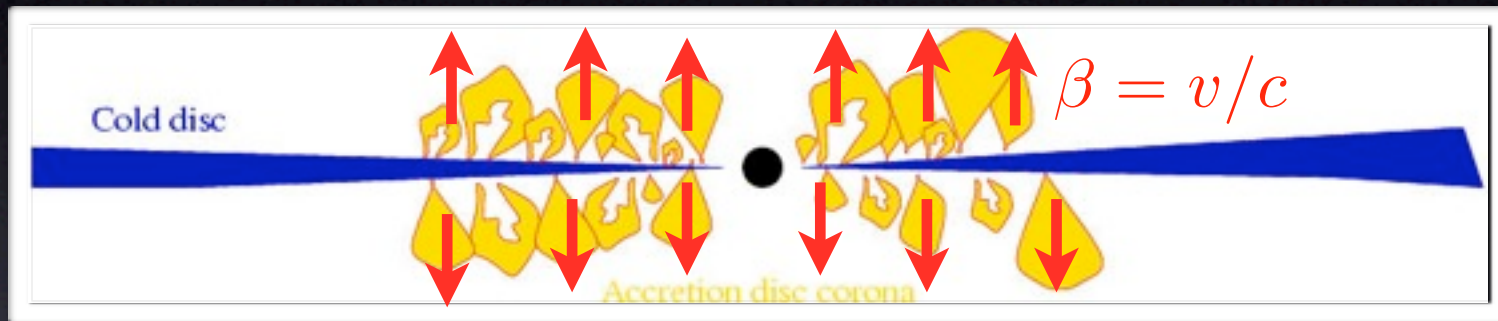
HIGH SOFT STATE:

disc blackbody and reflection: strong / Corona: NON-THERMAL (or hybrid) Comptonisation



Alternative models for the hard state

- Accretion disc corona outflowing with mildly relativistic velocity above a cold (i.e. non-radiating) thin disc



(Beloborodov 1999; Malzac Beloborodov & Poutanen 2001)

- Hard state associated with radio jet: X-ray Jet Models

(Markoff et al. 2001, 2005; Reig et al. 2003; Giannios et al. 2004; Kylafis et al. 2008)

BELM: a code to model radiation and kinetic processes in the corona

➡ Evolution of electrons and photon energy distributions in a fully ionised, magnetised plasma (radiation, acceleration and Coulomb processes)

Belmont, Malzac & Marcowith 2008

see also Vurm & Poutanen 2009 and talk by J. Poutanen

BELM: a code to model radiation and kinetic processes in the corona

➡ Evolution of electrons and photon energy distributions in a fully ionised, magnetised plasma (radiation, acceleration and Coulomb processes)

Belmont, Malzac & Marcowith 2008

see also Vurm & Poutanen 2009 and talk by J. Poutanen

MODEL:

- Continuous POWER-LAW electron acceleration in magnetised plasma
 - ➡ Cooling and thermalisation through synchrotron self-Compton + e-e Coulomb
 - ➡ Equilibrium distribution: Maxwellian+ non-thermal tail

BELM: a code to model radiation and kinetic processes in the corona

➡ Evolution of electrons and photon energy distributions in a fully ionised, magnetised plasma (radiation, acceleration and Coulomb processes)

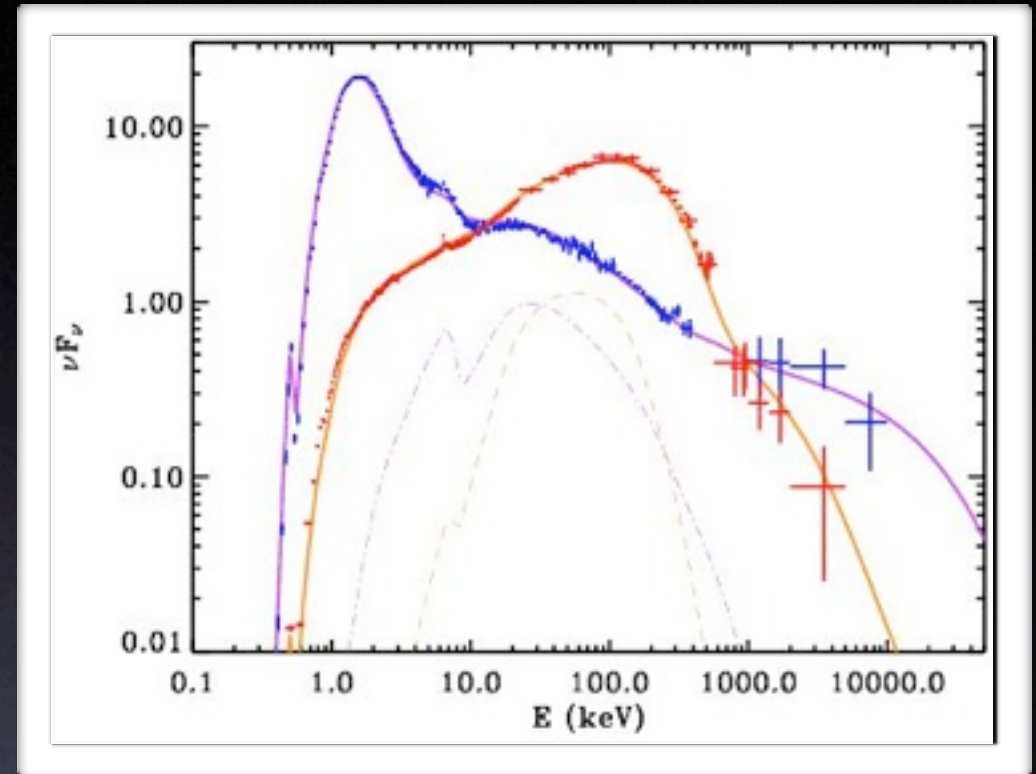
Belmont, Malzac & Marcowith 2008

see also Vurm & Poutanen 2009 and talk by J. Poutanen

MODEL:

- Continuous POWER-LAW electron acceleration in magnetised plasma
 - ➡ Cooling and thermalisation through synchrotron self-Compton + e-e Coulomb
 - ➡ Equilibrium distribution: Maxwellian+ non-thermal tail
- If corona illuminated by accretion disc: additional Compton cooling on external photons

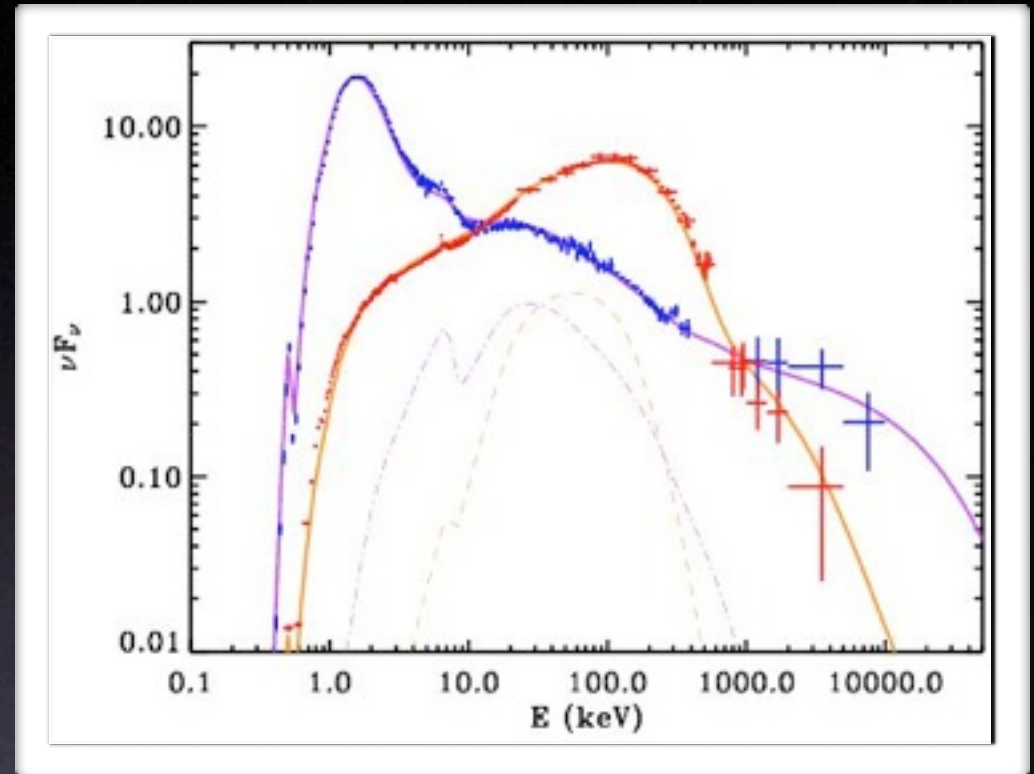
Comparisons to Cygnus X-I spectra



(Malzac & Belmont MNRAS 2009 ; Poutanen & Vurm 2009)

Comparisons to Cygnus X-I spectra

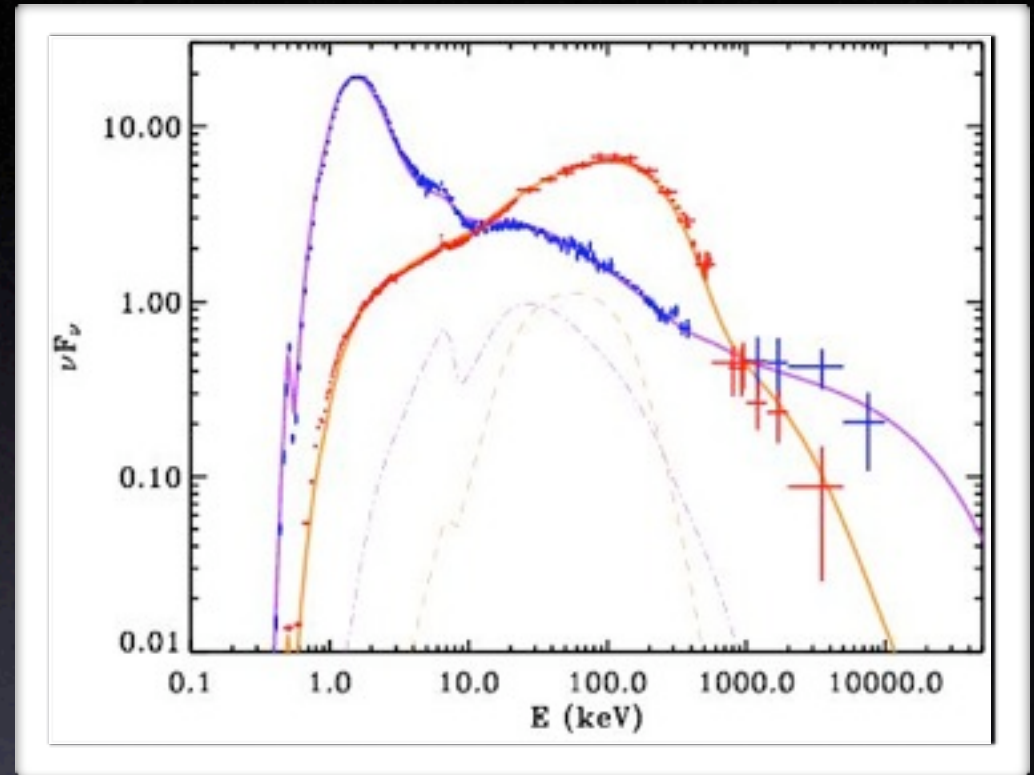
- Both states consistent with pure non-thermal acceleration models
- Different coronal temperatures due to more cooling by thermal disc photons in Soft state



(Malzac & Belmont MNRAS 2009 ; Poutanen & Vurm 2009)

Comparisons to Cygnus X-I spectra

- Both states consistent with pure non-thermal acceleration models
- Different coronal temperatures due to more cooling by thermal disc photons in Soft state

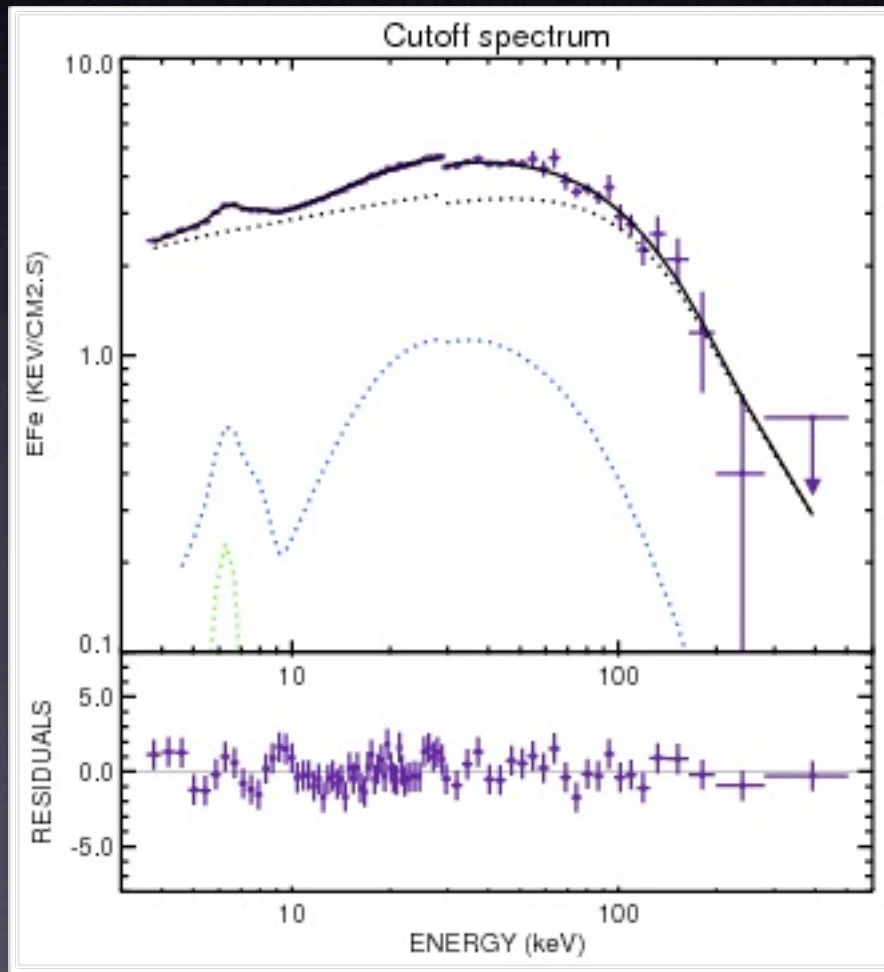


- Magnetic field in hard state: $U_B/U_R < 0.3$
➔ corona unlikely to be powered by magnetic field

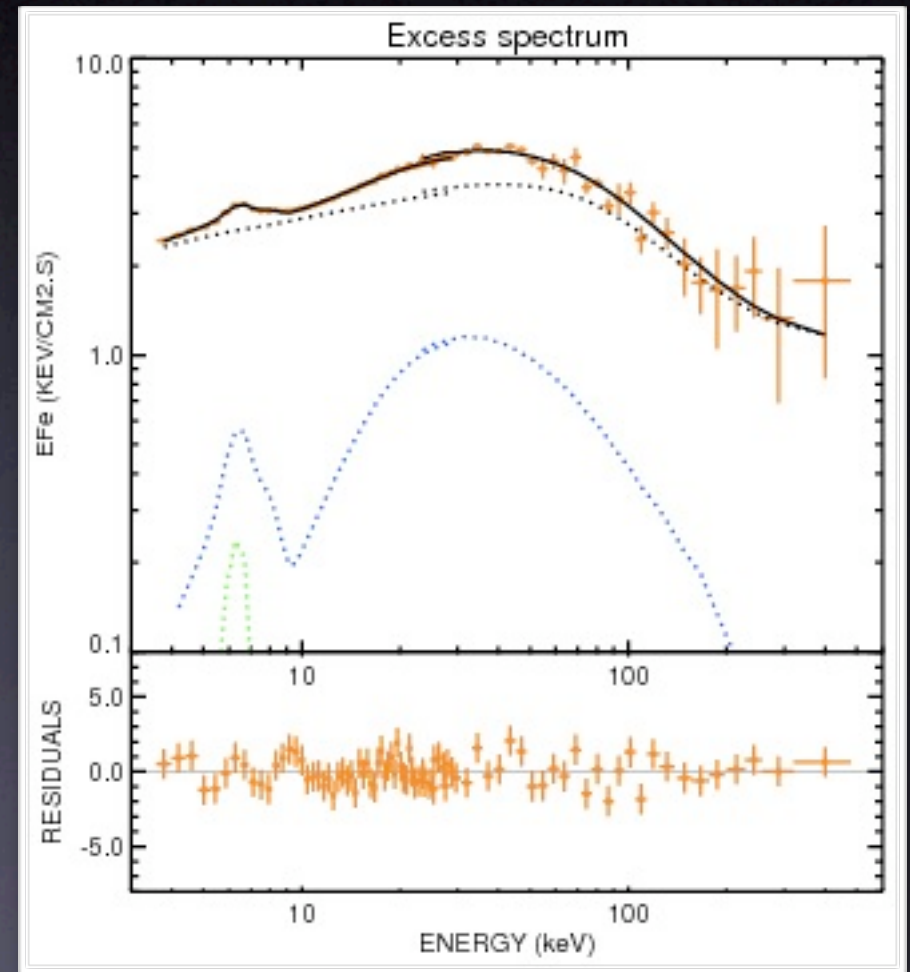
(Malzac & Belmont MNRAS 2009 ; Poutanen & Vurm 2009)

RXTE/INTEGRAL observation of GX 339-4 during a bright Low Hard State

- Spectral variability detected above 100 keV on time-scales of a few hours




lower flux: spectrum with high energy cut-off



higher fluxes: spectrum with high energy excess

Spectral fits with BELM

Spectrum	Γ_{inj}	γ_{max}	l_B	τ_{ion}	τ_T	kT_e (keV)	$\Omega/2\pi$	ξ	$\chi^2_v(\text{dof})$
cutoff	$2.5f$	$15.2^{+5.5}_{-1.3}$	283^{+43}_{-36}	$2.30^{+0.17}_{-0.16}$	$2.30^{+0.17}_{-0.16}$	33.4	$0.43^{+0.04}_{-0.03}$	301^{+123}_{-100}	0.96(68)
excess	$2.5f$	187^{+111}_{-105}	$27.3^{+4.0}_{-3.4}$	$2.49^{+0.18}_{-0.18}$	$2.73^{+0.18}_{-0.18}$	23.9	$0.39^{+0.04}_{-0.04}$	310^{+155}_{-126}	0.92(74)

-  LHS of GX339-4 is consistent with pure non-thermal injection SSC model, B could be larger than in Cyg X-1 ($B \leq 10^7$ gauss)

(Droulans, Belmont, Malzac & Jourdain, in preparation)

Spectral fits with BELM

Spectrum	Γ_{inj}	γ_{max}	l_B	τ_{ion}	τ_T	kT_e (keV)	$\Omega/2\pi$	ξ	$\chi^2_v(dof)$
cutoff	$2.5f$	$15.2^{+5.5}_{-1.3}$	283^{+43}_{-36}	$2.30^{+0.17}_{-0.16}$	$2.30^{+0.17}_{-0.16}$	33.4	$0.43^{+0.04}_{-0.03}$	301^{+123}_{-100}	0.96(68)
excess	$2.5f$	187^{+111}_{-105}	$27.3^{+4.0}_{-3.4}$	$2.49^{+0.18}$	$2.73^{+0.18}$	23.9	$0.39^{+0.04}_{-0.04}$	310^{+155}_{-126}	0.92(74)

- LHS of GX339-4 is consistent with pure non-thermal injection SSC model, B could be larger than in Cyg X-1 ($B \leq 10^7$ gauss)
- Variability of high energy tail can be interpreted in terms of fluctuation of magnetic field and γ_{max}
 - ➔ B changes by approximately a factor of 3
 - ➔ The fitted dependence between γ_{max} and B is: $\gamma_{max} \propto 1/B^2$

(Droulans, Belmont, Malzac & Jourdain, in preparation)

Model with hot protons

(Malzac & Belmont MNRAS 2009)

Model with hot protons

In addition to non-thermal acceleration we now assume that electrons are heated through Coulomb interactions with a population of hot thermal protons (two-temperature plasma):

(Malzac & Belmont MNRAS 2009)

Model with hot protons

In addition to non-thermal acceleration we now assume that electrons are heated through Coulomb interactions with a population of hot thermal protons (two-temperature plasma):

- Good agreement with data
but non-thermal electron injection is required in both Low Hard and High Soft states

(Malzac & Belmont MNRAS 2009)

Model with hot protons

In addition to non-thermal acceleration we now assume that electrons are heated through Coulomb interactions with a population of hot thermal protons (two-temperature plasma):

- Good agreement with data
but non-thermal electron injection is required in both Low Hard and High Soft states
- Temperature of hot protons in hard state:
 $T_i < 2 \times 10^{10} \text{ K}$ or $T_i/T_e < 10$
➔ proton temperature much lower than standard two-temperature accretion disc solutions

(Malzac & Belmont MNRAS 2009)

Jet energetics of Cygnus X-1

🌍 Jet powered nebula: $P_j \simeq L_X \simeq 2 \times 10^{37} \text{ erg s}^{-1}$ (Gallo et al. 2005, Russell et al. 2007)



© 2009 Stephen G. Cullen/LightBuckets

Jet energetics of Cygnus X-1

🌍 Jet powered nebula: $P_j \simeq L_X \simeq 2 \times 10^{37} \text{ erg s}^{-1}$ (Gallo et al. 2005, Russell et al. 2007)



Consequences for the accretion efficiency

- In soft state, accretion via thin disc is efficient.
Assuming no jet, the accretion efficiency:

$$\eta_s = L_s / \dot{M}_s c^2$$

Consequences for the accretion efficiency

- In soft state, accretion via thin disc is efficient.
Assuming no jet, the accretion efficiency:

$$\eta_s = L_s / \dot{M}_s c^2$$

- In hard state:

$$\eta_h = \frac{L_h + P_J}{(1 - f_j) \dot{M}_h c^2}$$

Consequences for the accretion efficiency

- In soft state, accretion via thin disc is efficient.
Assuming no jet, the accretion efficiency:

$$\eta_s = L_s / \dot{M}_s c^2$$

- In hard state:

$$\eta_h = \frac{L_h + P_J}{(1 - f_j) \dot{M}_h c^2}$$

- Observations: $P_J \simeq L_h$, $\frac{L_s}{L_h} \leq 4$, and $\dot{M}_s \geq \dot{M}_h$

Consequences for the accretion efficiency

- In soft state, accretion via thin disc is efficient.
Assuming no jet, the accretion efficiency:

$$\eta_s = L_s / \dot{M}_s c^2$$

- In hard state:

$$\eta_h = \frac{L_h + P_J}{(1 - f_j) \dot{M}_h c^2}$$

- Observations: $P_J \simeq L_h$, $\frac{L_s}{L_h} \leq 4$, and $\dot{M}_s \geq \dot{M}_h$

$$\Rightarrow \eta_h \geq \eta_s / 2$$

- ➡ accretion proceeds efficiently in the hard state
- ➡ cannot be strongly advection dominated

Constraints on jet terminal velocity

$$P_J = f_j \dot{M}_h (\gamma_\infty - 1) c^2$$

(Malzac, Belmont & Fabian, MNRAS, 2009)

Constraints on jet terminal velocity

$$P_J = f_j \dot{M}_h (\gamma_\infty - 1) c^2$$

$$P_J + L_h = \eta_h (1 - f_j) \dot{M}_h c^2$$

(Malzac, Belmont & Fabian, MNRAS, 2009)

Constraints on jet terminal velocity

$$P_J = f_j \dot{M}_h (\gamma_\infty - 1) c^2$$

$$P_J + L_h = \eta_h (1 - f_j) \dot{M}_h c^2$$

$$L_s = \eta_s \dot{M}_s c^2$$

(Malzac, Belmont & Fabian, MNRAS, 2009)

Constraints on jet terminal velocity

$$P_J = f_j \dot{M}_h (\gamma_\infty - 1) c^2$$

$$P_J + L_h = \eta_h (1 - f_j) \dot{M}_h c^2 \quad \Rightarrow \quad \gamma_\infty = 1 + \frac{\eta_s j}{\lambda(1 + j) \frac{\eta_s}{\eta_h}}$$

$$L_s = \eta_s \dot{M}_s c^2$$

(Malzac, Belmont & Fabian, MNRAS, 2009)

Constraints on jet terminal velocity

$$P_J = f_j \dot{M}_h (\gamma_\infty - 1) c^2$$

$$P_J + L_h = \eta_h (1 - f_j) \dot{M}_h c^2 \quad \Rightarrow \quad \gamma_\infty = 1 + \frac{\eta_s j}{\lambda (1 + j) \frac{\eta_s}{\eta_h}}$$

$$L_s = \eta_s \dot{M}_s c^2$$

$$\lambda = \frac{L_s \dot{M}_h}{L_h \dot{M}_s} < 4$$

$$j = P_J / L_h \simeq 1$$

(Malzac, Belmont & Fabian, MNRAS, 2009)

Constraints on jet terminal velocity

$$P_J = f_j \dot{M}_h (\gamma_\infty - 1) c^2$$

$$P_J + L_h = \eta_h (1 - f_j) \dot{M}_h c^2 \quad \Rightarrow \quad \gamma_\infty = 1 + \frac{\eta_s j}{\lambda (1 + j) \frac{\eta_s}{\eta_h}}$$

$$L_s = \eta_s \dot{M}_s c^2$$

$$\lambda = \frac{L_s \dot{M}_h}{L_h \dot{M}_s} < 4$$

$$j = P_J / L_h \simeq 1$$

$$\Rightarrow \beta_\infty > 0.1$$

$$\text{and most likely: } 0.8 \gtrsim \beta_\infty \gtrsim 0.3$$

(Malzac, Belmont & Fabian, MNRAS, 2009)

Is the X-ray emission from the jet ?

● Mass flux conservation along the jet implies:

$$\dot{M}_J = \pi r_0^2 n_0 \beta_0 \gamma_0 m_p c = \frac{P_J}{(\gamma_\infty - 1)c^2}$$

- jet power P_J ($\simeq 2 \times 10^{37}$ erg/s)
- terminal velocity β_∞ (> 0.1)
- jet density at base n_0
- jet radius at base r_0
- initial jet velocity β_0

(Malzac, Belmont & Fabian, MNRAS, 2009)

Is the X-ray emission from the jet ?

- Mass flux conservation along the jet implies:

$$\dot{M}_J = \pi r_0^2 n_0 \beta_0 \gamma_0 m_p c = \frac{P_J}{(\gamma_\infty - 1)c^2}$$

- jet power P_J ($\simeq 2 \times 10^{37}$ erg/s)
- terminal velocity β_∞ (> 0.1)
- jet density at base n_0
- jet radius at base r_0
- initial jet velocity β_0

- If X-ray emission is Comptonisation in the jet:

$$\tau_T = n_0 \sigma_T r_0 > 1$$

(Malzac, Belmont & Fabian, MNRAS, 2009)

Is the X-ray emission from the jet ?

- Mass flux conservation along the jet implies:

$$\dot{M}_J = \pi r_0^2 n_0 \beta_0 \gamma_0 m_p c = \frac{P_J}{(\gamma_\infty - 1)c^2}$$

- jet power P_J ($\simeq 2 \times 10^{37}$ erg/s)
- terminal velocity β_∞ (> 0.1)
- jet density at base n_0
- jet radius at base r_0
- initial jet velocity β_0

- If X-ray emission is Comptonisation in the jet:

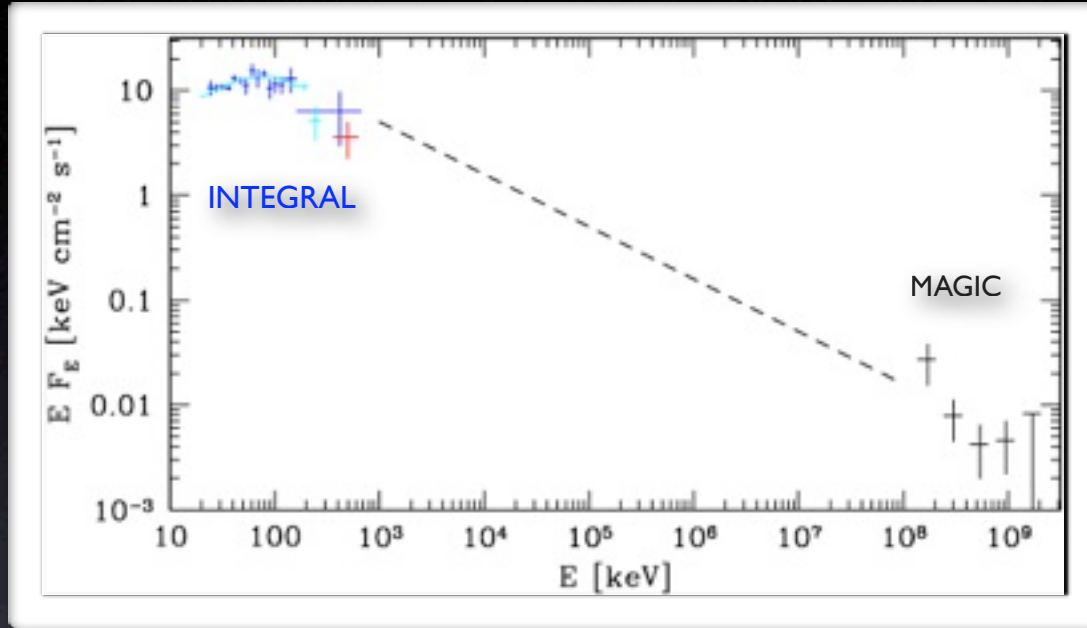
$$\tau_T = n_0 \sigma_T r_0 > 1$$

$$\Rightarrow \beta_0 < 0.1$$

➡ velocity of X-ray emitting region much slower than
in any X-ray jet model

(Malzac, Belmont & Fabian, MNRAS, 2009)

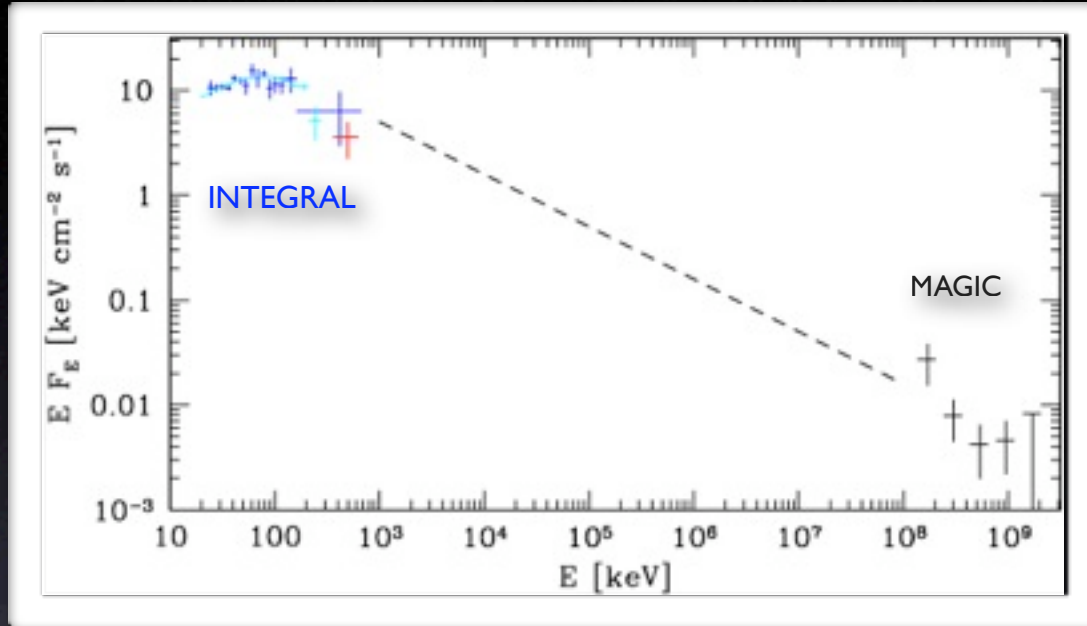
TeV detection of Cyg X-1 by MAGIC (Albert et al., 2007)



During MAGIC detection Cyg X-1 was in a stable hard state. The brightest ever observed. Otherwise nothing unusual in X-ray light curves or spectra.

(Malzac et al. A&A 2008)

TeV detection of Cyg X-1 by MAGIC (Albert et al., 2007)



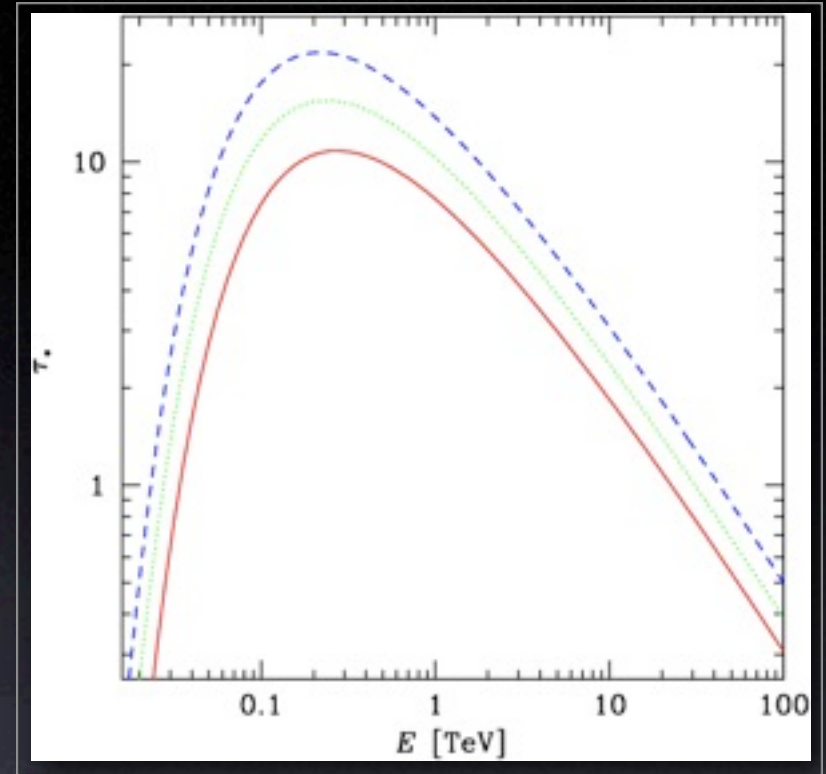
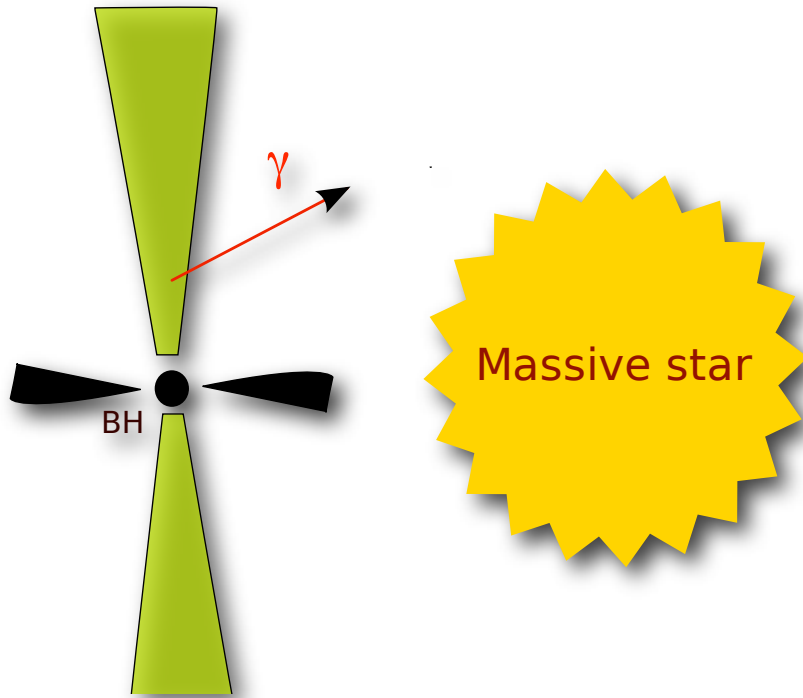
During MAGIC detection Cyg X-1 was in a stable hard state. The brightest ever observed. Otherwise nothing unusual in X-ray light curves or spectra.

(Malzac et al. A&A 2008)

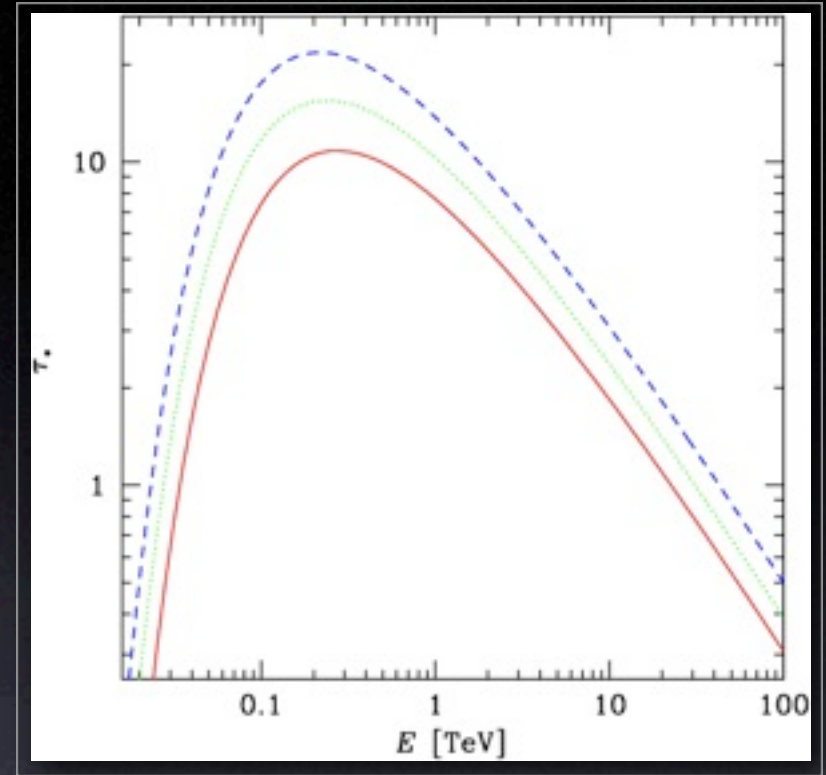
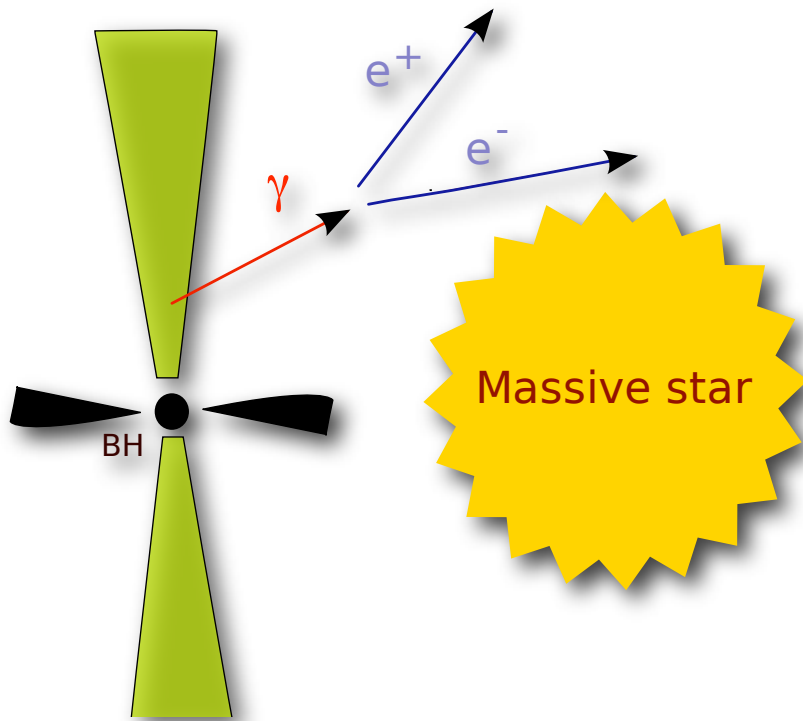
Origin of TeV emission ?

- ➡ in the X-ray corona radiative cooling is too strong to accelerate particles up to TeV energies
- ➡ shocks in jet, or interaction of outflow with wind of companion star
- ➡ relatively close to the black hole
- ➡ main constraint: pair absorption in the field of companion star

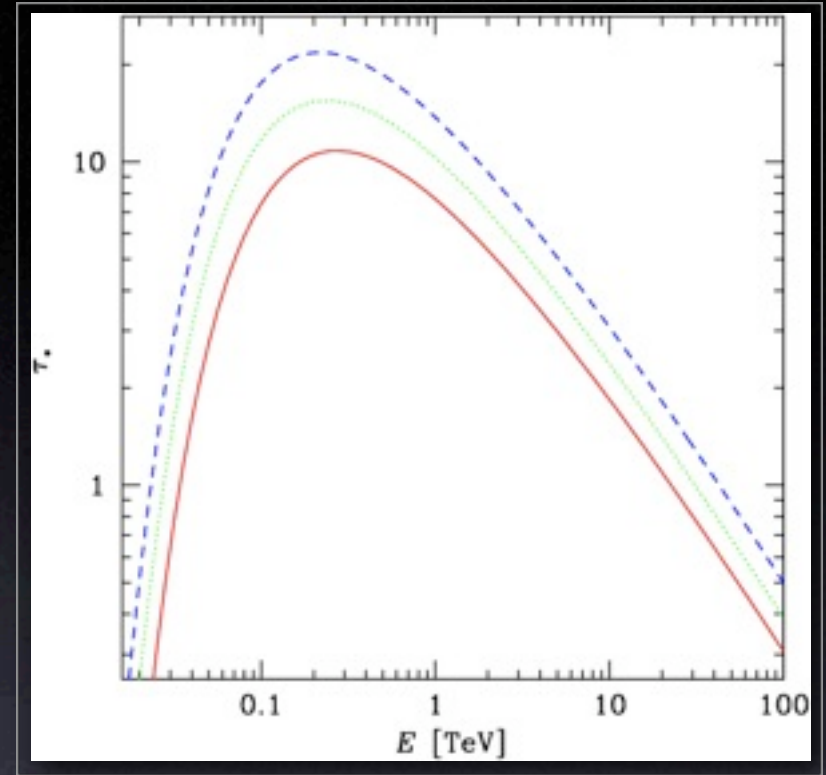
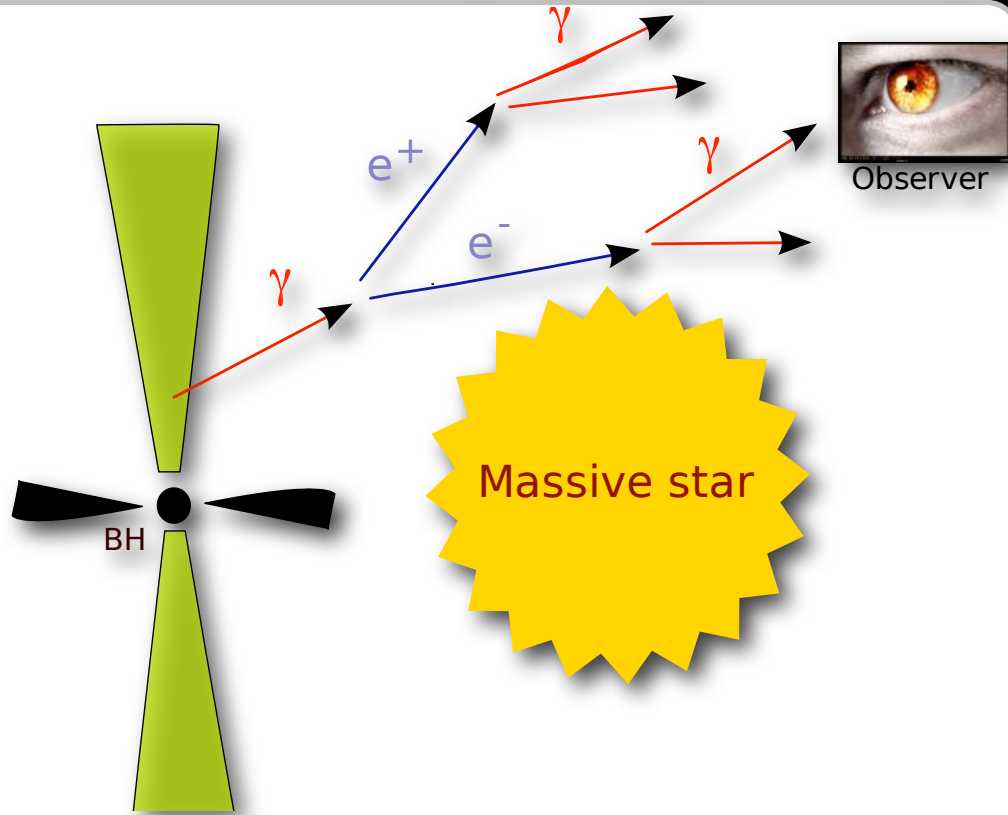
➡ BH behind companion: pair absorption in radiation field of star



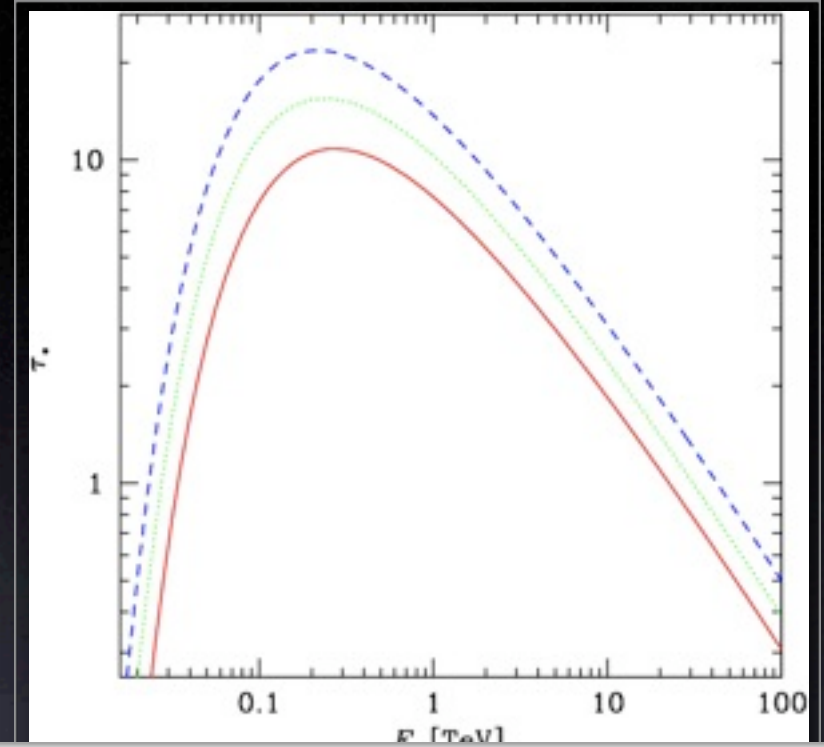
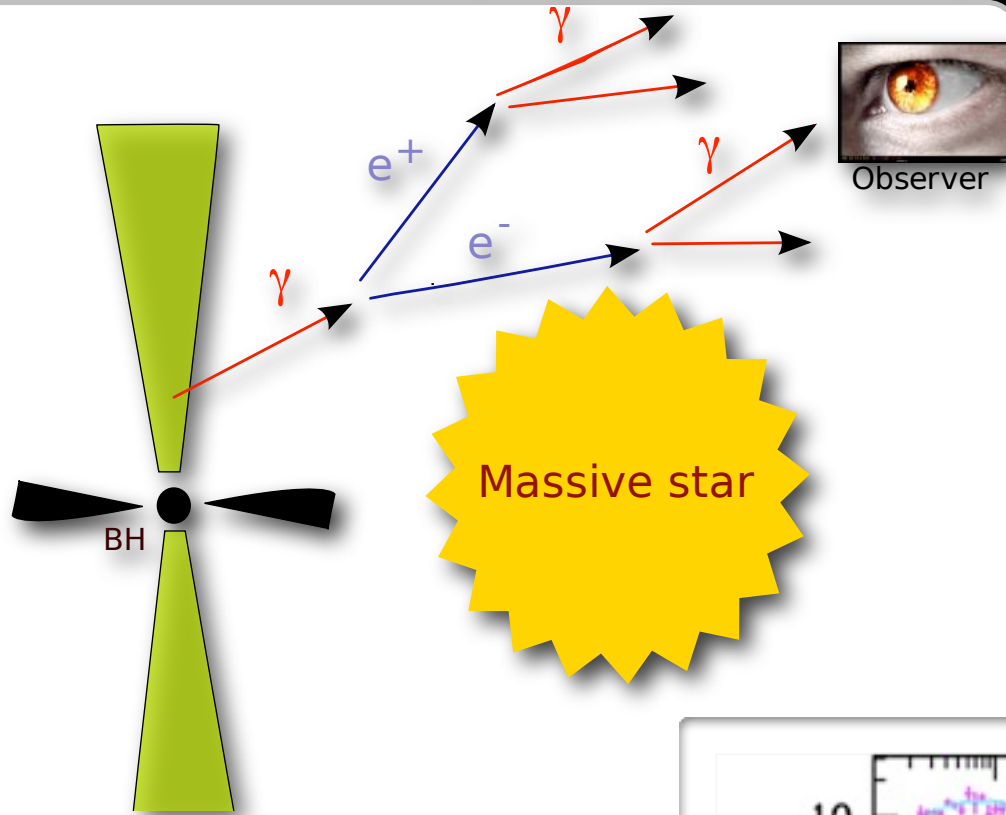
➡ BH behind companion: pair absorption in radiation field of star



➡ BH behind companion: pair absorption in radiation field of star

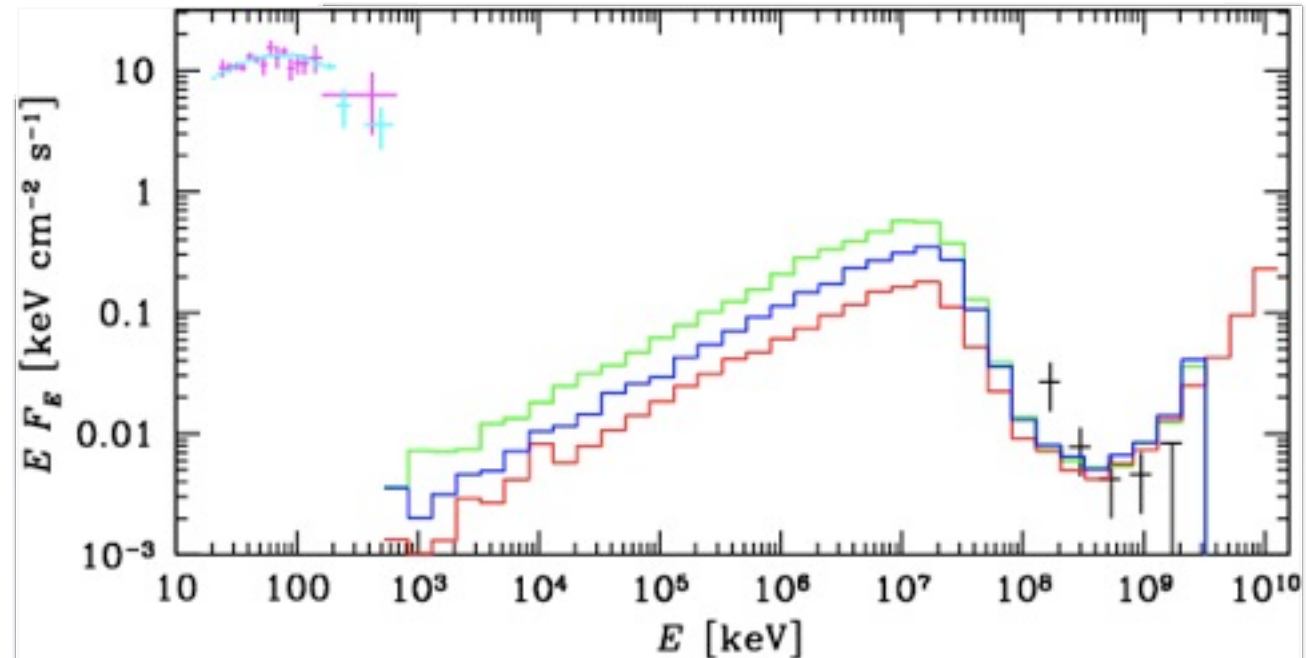


➔ BH behind companion: pair absorption in radiation field of star



➔ spatially extended pair cascades

(Zdziarski, Malzac, Bednarek, MNRAS, 2009)



Conclusions:

- Radio jet velocity appears to be mildly relativistic
- X-ray Jet models: appear to be ruled out
(although jet may contribute to gamma-ray emission)
- Outflowing accretion disc corona:
 - corona not powered through magnetic dissipation
 - small mass loading on the jet
- Hot accretion flows models:
 - appear quite efficient
 - temperature of electrons and protons are comparable